Optics for Engineers Week 5

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Week 5 Agenda

- Coherent and Incoherent Light
- Mach–Zehnder Interferometer
- Michaelson Interferometer and Optical Testing
- Doppler Lidar
- Fabry–Perot Interfereometer
- Laser Cavities
- Dielectric Coatings

Superposition



- Output Toward the Right (Similar for Downward)
- Coherent Addition

$$E = E_1 + E_2$$
 $I = |E|^2 = EE^*$

• Incoherent Addition

$$I = I_1 + I_2$$

Coherent Superposition



- Inputs Both Have I = 1
- Outputs Both Have $0 \le E_1 + E_2 \le 2$
- But $I_{right} + I_{down} = I_{from top} + I_{from left}$

Measuring the Field Amplitude Is Hard

- Easy for Ocean Wave Height
- Easy for Acoustic Pressure
- Even Easy for Radio Waves
- No Direct Measurement for Light
 - Terahertz Frequencies
 - Sub–Micrometer
 Wavelengths
- Use Interferometry
 - Mix With Known
 Reference Wave
 - Measure Irradiance
 - Variations in Space or Time



Interferometry Equations

• Irradiance

$$I = \frac{\left(E_{1}^{*} + E_{2}^{*}\right)\left(E_{1} + E_{2}\right)}{Z}$$

• Expand

$$I = \frac{E_1^* E_1 + E_2^* E_2 + E_1^* E_2 + E_1 E_2^*}{Z}$$

- First Two Terms are "DC" Terms
- Third and Fourth are "Mixing" Terms
- Complex Conjugate Pair (Real Sum)

$$I_{mix} = \frac{E_1 E_2^*}{Z}$$
 and $I_{mix}^* = \frac{E_1^* E_2}{Z}$

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Mixing Terms

• Complex Conjugates Add to Real Value

$$I_{mix} = \frac{E_1 E_2^*}{Z}$$
 and $I_{mix}^* = \frac{E_1^* E_2}{Z}$

• Magnitude

$$|I_{mix}| = |I_{mix}^*| = \sqrt{I_1 I_2}$$

• Total Irradiance

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\phi_2 - \phi_1)$$

• Random Fields: Incoherent Superposition

$$\bar{I}_{mix} = 0 \qquad \bar{I} = I_1 + I_2$$

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Mach–Zehnder Interferometer



$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\phi_2 - \phi_1)$$

• Add Gas Pressure to Cell; $n \uparrow$, $\Delta = \Delta OPL \uparrow$

$$\Delta = \delta \left(n \ell_c \right)$$

$$\delta\phi_1 = k\Delta = 2\pi \frac{\Delta}{\lambda} = 2\pi \frac{\ell_c}{\lambda} \delta n$$

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Fringe Amplitude and Contrast

• Total Signal

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$$I = I_0 \left(R_1 T_2 + T_1 R_2 + 2\sqrt{R_1 T_2 T_1 R_2} \cos \delta \phi \right)$$

• Fringe Amplitude

$$I_m = I_{max} - I_{min}$$

$$\sqrt{R_1 T_1 R_2 T_2} I_0 = 0.25 I_0$$
 for $R_1 = R_2 = 0.5$

• Fringe Contrast Defined

$$V = \frac{I_{max} - I_{min}}{I_{max} + I_{min}} \qquad (0 \le V \le 1) \qquad V = 2\frac{\sqrt{R_1 T_2 T_1 R_2}}{R_1 T_2 + T_1 R_2}$$

• For $R_1 = R_2 = R$ and $(T_1 = T_2 = T)$

$$I = I_0 2RT \left(1 + \cos \delta \phi\right)$$

Michelson Interferometer



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Optical Testing

Synthetic Results Illustrating Fringe Patterns



Perfect Surface 2–Wavelength Bump 0.2–Wavelength Bump (4 Fringes) (0.4 Fringes)

Coherent Laser Radar (Lidar)

- Similar to Mach–Zehnder (Modified Mach–Zehnder?)
- Common Transmit/Receive Aperture: Use T/R Switch (Ch. 6)
 - Transmitter Polarization: *P* at Beamsplitters
 - Receiver Polarization: S: Need HWP in Reference (LO)
 - Ideally QWP Between Telescope and Target to Reduce Narcissus (Not Practical)
- BS1 and Recombining Beamsplitter High Reflectivity

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Doppler Velocity

- Target: Dust, Fo, Rain, Snow, Smoke, etc.
- Doppler Equation from Source to Target

 $2\pi f_d = \mathbf{k} \cdot \mathbf{v}$ (> 0 for Approaching Velocities)

$$f_d = \frac{v_{parallel}}{\lambda}$$
 (Moving Source or Detector)

• Doppler Lidar (or Radar) on Round Trip

$$f_{DR} = 2\frac{v_{parallel}}{\lambda}$$

 $(f_{DR} = 100$ kHz for $v_{parallel} = 0.54$ m/s at $\lambda = 10.59$ μ m)

Pulsed Laser Radar

• Range Resolution and Velocity Resolution

$$\delta r = rac{c au}{2} \qquad \delta f_{DR} pprox rac{1}{ au} \qquad \delta v_{parallel} pprox rac{\lambda}{2 au}$$

- \bullet Hundreds of Meters and m/s Typical at 10.59 $\mu{\rm m}$
- Average Power

$$P_{avg} = P_{laser} \times \tau \times PRF$$

• Range Ambiguity May Limit PRF



Fabry-Perot



 Resonant Frequencies and Free Spectral Range

 $f = N f_0$

$$FSR = f_0 = \frac{c}{2\ell}$$

• Recirculating Power

$$P_{recirculating} = \frac{P_{out}}{T_2} =$$

$$\frac{P_{out}}{1 - R_2} = \frac{P_0}{1 - R_2}$$

$$P_N = (R_1 R_2)^N$$

 $N = -\log 2/\log \left(R_1 R_2\right)$

e.g.
$$R_1 = R_2 = 0.999$$
:

N = 346

• Resolution of a Longer Interferometer

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Fabry–Perot Equations



The Infinite Sum, or "Barber's Chair" Approach



Computing the Sum

$$E_t = E_0 \tau_1 \tau_2 e^{jk\ell} \frac{1}{1 - \rho_1' \rho_2 e^{2jk\ell}}$$

$$E_r = E_0 \rho_1 + \tau_1 \tau_1' e^{2jk\ell} \frac{1}{1 - \rho_1' \rho_2 e^{2jk\ell}}$$

$$T = \left| \frac{E_t}{E_0} \right|^2 \qquad T = \tau_1 \tau_2 \tau_1^* \tau_2^* \frac{1}{1 - \rho_1' \rho_2 e^{2jk\ell}} \frac{1}{1 - (\rho')_1^* \rho_2^* e^{-2jk\ell}}$$
$$T = T_1 T_2 \frac{1}{1 - 2\sqrt{R_1 R_2} \cos\left(2k\ell\right) + R_1 R_2}$$

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Fabry–Perot Transmission



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Thin Films: Approach



- Normal Incidence
- Wave from Left

$$E_A = E_i + E_r \qquad E_D = E_t$$

$$H_A = H_i - H_r \qquad H_D = H_t$$

• Electric Field BC

$$E_B = E_A \qquad E_C = E_D$$

- Magnetic Field BC
 - $H_B = H_A \qquad H_C = H_D$
 - $\frac{E_B}{nZ_0} = \frac{E_A}{n_0Z_0} \qquad \frac{E_c}{nZ_0} = \frac{E_D}{n_tZ_0}$



 $\frac{E_B}{=} = \frac{E_A}{=} \frac{E_C}{=} \frac{E_D}{=}$ n_0 n n_t n

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Thin Films: In the Medium

• Right–Propagating

 $E_{right}e^{inkz}$ and $H_{right}e^{inkz} = \frac{1}{nZ_0}E_{right}e^{inkz}$ • Left–Propagating $E_{left}e^{-inkz}$ and $H_{left}e^{-inkz} = -\frac{1}{nZ_0}E_{left}e^{-inkz}$

• Boundaries

$$E_B = E_{left} + E_{right}$$
 $E_c = E_{left}e^{-jnk\ell} + E_{right}e^{jnk\ell}$

• Result

$$E_i + E_r = E_t \cos(nk\ell) - E_t \frac{n_t}{n} \sin(nk\ell)$$

$$n_0 E_i - n_0 E_r = -inE_t \sin(nk\ell) + n_t E_t \cos(nk\ell)$$

Thin Films: Matrix Equation

• Previous Result

$$E_i + E_r = E_t \cos(nk\ell) - E_t \frac{n_t}{n} \sin(nk\ell)$$

 $n_0 E_i - n_0 E_r = -inE_t \sin(nk\ell) + n_t E_t \cos(nk\ell)$

• Matrix Equation

$$\begin{pmatrix} 1\\n_0 \end{pmatrix} E_i + \begin{pmatrix} 1\\-n_0 \end{pmatrix} E_r = \begin{pmatrix} \cos(nk\ell) & -\frac{i}{n}\sin(nk\ell)\\-in\sin(nk\ell) & \cos(nk\ell) \end{pmatrix} \begin{pmatrix} 1\\n_t \end{pmatrix} E_t$$

• Characteristic Matrix

$$\mathcal{M} = \begin{pmatrix} \cos(nk\ell) & -\frac{i}{n}\sin(nk\ell) \\ -in\sin(nk\ell) & \cos(nk\ell) \end{pmatrix}$$

$$\begin{pmatrix} 1\\n_0 \end{pmatrix} + \begin{pmatrix} 1\\-n_0 \end{pmatrix} \rho = \mathcal{M} \begin{pmatrix} 1\\n_t \end{pmatrix} \tau$$

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Multiple Layers

$$\mathcal{M} = \begin{pmatrix} \cos(nk\ell) & -\frac{i}{n}\sin(nk\ell) \\ -in\sin(nk\ell) & \cos(nk\ell) \end{pmatrix} \begin{pmatrix} 1 \\ n_0 \end{pmatrix} + \begin{pmatrix} 1 \\ -n_0 \end{pmatrix} \rho = \mathcal{M} \begin{pmatrix} 1 \\ n_t \end{pmatrix} r$$

$$\mathcal{M} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

$$\mathcal{M} = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

$$\rho = \frac{An_0 + Bn_tn_0 - C - Dn_t}{An_0 + Bn_tn_0 + C + Dn_t}$$

$$\mathcal{M} = \mathcal{M}_1 \mathcal{M}_2 \mathcal{M}_3 \dots$$

$$\tau = \frac{2n_0}{An_0 + Bn_tn_0 + C + Dn_t}$$

Dielectric Stacks

- High Reflectivity (Often Better than Metal)
- Anti-Reflection Coatings or Stacks
- Narrow–Band Filters, Mirrors, etc.
- Bandpass Devices that Are Not Narrow–Band
- Hot Mirror or Cold Mirror
- Long–Pass Dichroic
- Short–Pass Dichroic
- Beamsplitters (Specific Reflectivity, Angle, Polarization, Wavelength Range, *etc.*)

High–Reflectance Stack (1)



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High–Reflectance Stack (2)

• One Pair of $\lambda/4$ Layers

$$\mathcal{M}_p = \begin{pmatrix} -n_\ell/n_h & 0\\ 0 & -n_h/n_\ell \end{pmatrix}$$

• Multiple Pairs

$$\mathcal{M}_N = \begin{pmatrix} (-n_\ell/n_h)^N & 0\\ 0 & (-n_h/n_\ell)^N \end{pmatrix}$$

• Reflectivity

$$R = \left(\frac{\left(\frac{n_{\ell}}{n_{h}}\right)^{2N} - \frac{n_{t}}{n_{0}}}{\left(\frac{n_{\ell}}{n_{h}}\right)^{2N} + \frac{n_{t}}{n_{0}}}\right)^{2}$$

- Narrow Band $(\lambda/4)$
- Almost Indep.of n_t , n_0

- Example
 - Zinc Sulfide, $n_h = 2.3$
 - Magnesium Fluoride,
 - $n_{\ell} = 1.35$
 - 8 Layers

$$N = 4 \rightarrow R = 0.97$$

- 30 Layers
 - $N = 15 \rightarrow R = 0.999$
- For Lasers, eg. HeNe
- Compare Metal $(\approx 0.96 \text{ Typical})$
 - Near-Zero Heating

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Anti–Reflection Coating



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Multi–Layer AR Stacks



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