## Optics for Engineers Chapter 10

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#### Coherent Sum

• Adding Fields (Chapter 7)

$$I = E_1^* E_1 + E_2^* E_2 + E_1^* E_2 + E_1 E_2^*$$

• Mach–Zehnder Interferometer

$$E_1(t) = E_s\left(t - \frac{z_1}{c}\right) \qquad E_2(t) = E_s\left(t - \frac{z_2}{c}\right)$$

• Mixing Term: (Source,  $E_s(t) = E_0 e^{j\omega t}$ , with  $k = 2\pi/\lambda = \omega/c$ )  $E_2 E_1^* = E_s(t - z_2/c) E_s^*(t - z_1/c) =$ 

$$E_0 (t - z_2/c) E_0^* (t - z_1/c) e^{j(\phi_2 - \phi_1)}$$

$$\phi_1 = k z_1 \qquad \phi_2 = k z_2$$

#### Irradiance of Coherent Sum

• Coherent Light as in Chapter 7

$$I = I_1 + I_2 + 2\sqrt{I_1 I_2} \cos(\phi_2 - \phi_1)$$

• Visibility for Equal Irradiances

$$V = \frac{I_{max} - I_{min}}{I_{max} + I_{min}} = \frac{4I_1 - 0}{4I + 0} = 1$$

• Random Source ( $E_0$  Varies Randomly)

$$\langle E_2 E_1^* \rangle = \langle E_0 (t - z_2/c) E_0^* (t - z_1/c) \rangle e^{j(\phi_2 - \phi_1)}$$

$$\langle E_2 E_1^* \rangle = \langle E_0(t) E_0^*(t-\tau) \rangle e^{j(\phi_2 - \phi_1)} \qquad \tau = \frac{z_2 - z_1}{c}$$

#### Coherence

• Previous Result

$$\langle E_2 E_1^* \rangle = \langle E_0(t) E_0^*(t-\tau) \rangle e^{j(\phi_2 - \phi_1)} \qquad \tau = \frac{z_2 - z_1}{c}$$

Temporal Autocorrelation Function, Γ

$$I = I_1 + I_2 + 2\Re \left[ \Gamma e^{j(\phi_2 - \phi_1)} \right] \qquad \Gamma = \langle E_0^* (t - \tau) E_0 (t) \rangle$$

• Visibility with  $I_1 = I_2$ 

$$V = \frac{I_{max} - I_{min}}{I_{max} + I_{min}} = \frac{2I_1 (1 + |\Gamma|) - 2I_1 (1 - |\Gamma|)}{2I_1 (1 + |\Gamma|) + 2I_1 (1 - |\Gamma|)}$$
$$V = \gamma = \frac{4|\Gamma|}{4I} = \frac{|\Gamma|}{I}$$

### Temporal Autocorrelation Function

• General Definition (See Previous Page)

$$\Gamma = \langle E_0^* \left( t - \tau \right) E_0 \left( t \right) \rangle \qquad \gamma \left( \tau \right) = \frac{\Gamma \left( \tau \right)}{\left| E_1 \right| \left| E_2^* \right|}$$

• Short Times: Equal (Or at Least Correlated) Fields  $\Gamma = \langle E_1^*(t) E_1(t) \rangle$ 

$$\gamma( au) 
ightarrow 1$$
 as  $au 
ightarrow 0$ 

• Long Time: Uncorrelated Fields

$$\Gamma = \langle E_1^* (t - \tau) E_1 (t) \rangle = \langle E_1^* (t - \tau) \rangle \langle E_1 (t) \rangle$$
$$\gamma(\tau) \to 0 \quad \text{as} \quad \tau \to \infty$$

## Coherent and Incoherent Addition

• Coherent Addition (*e.g.* Short Times)

$$E = E_1 + E_2$$
  $I = |E|^2$ 

• Incoherent Addition (*e.g.* Long Times)

 $\langle I \rangle = \langle I_1 \rangle + \langle I_2 \rangle$  (Incoherent Addition)

• Need to Quantify Times

#### Digression: Mode–Locked Lasers

• Sum of Longitudinal Modes: *e.g.* 25 Modes

$$E = \sum_{m=-12}^{12} E_m = \sum_{m=-12}^{12} e^{i2\pi [f_{center} + m \times FSR]t}$$

 Transform–Limited Pulses Carrier Frequency Reduced for Clarity



#### Pulse Dispersion

Mode-LockedGlass (Lens or Fiber)Pulse notLaser $n=n(f)=n(c/\lambda)$ Limited

- Mode-Locked Laser Passing Through Glass
- "Chirped Pulses:" e.g.  $\frac{dt}{dt} = -(1/3) \times 10^{-18} \text{sec/Hz}$
- Pulses Broadened in Time (Not Transform-Limited)



#### Free–Running Multi–Mode Laser

- "Unlocked" Laser
- Random Output with No Predictable Pulses
- Occasional "Hot Pulse"



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## Multi–Mode Laser: Incoherent Addition

- 201 Modes 5MHz apart with Random Phases
- Short–Time Coherence ( $\ll 1/Bandwidth$ )



## Temporal Coherence (1)

• Fourier Transforms

$$\tilde{E}(\omega) = \int_{-\infty}^{\infty} E(t) e^{-j\omega t} dt$$

$$E(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{E}(\omega) e^{j\omega t} d\omega$$

• Autocorrelation

$$\Gamma(\tau) = \langle E^*(t-\tau) E(t) \rangle$$

$$E(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{E}(\omega) e^{j\omega t} d\omega$$
$$E^*(t-\tau) =$$
$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{E}^*(\omega') e^{-j\omega'(t-\tau)} d\omega'$$

• Combine and Simplify

$$\Gamma(\tau) = \left\langle \frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{E}^*(\omega') e^{-j\omega'(t-\tau)} d\omega' \dots \right.$$
$$\frac{1}{2\pi} \int_{-\infty}^{\infty} \tilde{E}(\omega) e^{j\omega t} d\omega \right\rangle$$
$$\Gamma(\tau) = \left\langle \frac{1}{4\pi^2} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \dots \right.$$
$$\tilde{E}^*(\omega') \tilde{E}(\omega) e^{j\omega'\tau} e^{j(\omega-\omega')t} d\omega' d\omega \right\rangle$$

• Exchange Average and Integrals

$$\Gamma(\tau) = \frac{1}{4\pi^2} \int_{-\infty}^{\infty} \dots$$

$$\left[\int_{-\infty}^{\infty} \left\langle \tilde{E}^{*}\left(\omega'\right) \tilde{E}\left(\omega\right) \right\rangle e^{j\omega'\tau} e^{j(\omega-\omega')t} d\omega' \right] d\omega$$

## Temporal Coherence (2)

• Previous Result

$$\Gamma(\tau) = \frac{1}{4\pi^2} \int_{-\infty}^{\infty} \left[ \int_{-\infty}^{\infty} \left\langle \tilde{E}^*(\omega') \tilde{E}(\omega) \right\rangle e^{j\omega'\tau} e^{j(\omega-\omega')t} d\omega' \right] d\omega$$

• Independent Sources with  $\left< \tilde{E}(\omega) \right> = 0$ 

 $\langle E_m E_n^* \rangle = I_m \delta_{m-n}$  (Discrete Spectrum; *e.g.* Multi-Mode Laser)  $\langle \tilde{E}^* (\omega') \tilde{E} (\omega) \rangle d\omega' = \tilde{I} (\omega') \delta (\omega - \omega') d\omega'$  (Continuous Spectrum *e.g.* LED)

• Result

$$\left[\int_{-\infty}^{\infty} \left\langle \tilde{E}^{*}\left(\omega'\right) \tilde{E}\left(\omega\right) \right\rangle e^{j\omega'\tau} e^{j\left(\omega-\omega'\right)t} d\omega' \right] = \tilde{I}\left(\omega\right) e^{j\omega\tau}$$

$$\Gamma(\tau) = \frac{1}{4\pi^2} \int_{-\infty}^{\infty} \tilde{I}(\omega) e^{j\omega\tau} d\omega \qquad \text{(Weiner-Khintchine)}$$

#### Coherence in Interfermetry

- Two Fields Shown (Solid and Dashed)
- Same Source with Different Time Delays
- Linewidth  $\approx 10 \text{GHz}$



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#### Example: LED

• Wavelength Spectrum (Gaussian)

$$\lambda_0 = 635$$
nm  $\lambda_{FWHM} = 20$ nm  $\left|\frac{\delta\lambda}{\lambda}\right| = \left|\frac{\delta f}{f}\right|$ 

• Gaussian Fourier Transform Pair

$$\gamma(\tau) = e^{-\left(t/t_{1/e}\right)^2} \qquad \frac{2\sqrt{\pi}}{\omega_{1/e}} \tilde{\gamma}(\omega) = e^{-\left(\omega/\omega_{1/e}\right)^2}$$

• Convert Widths and Complete Calculation

#### Gaussian Widths Defined



#### LED Example Conclusion

$$t_{1/2} = \sqrt{\ln 2} \frac{2\sqrt{2}}{\omega_{1/e}} = \ln 2 \frac{2\sqrt{2}}{2\pi f_{1/2}} \approx \frac{0.31}{f_{1/2}} = \frac{0.62}{f_{FWHM}}$$
$$z_{1/2} = ct_{1/2}$$
$$\frac{z_{1/2}}{\lambda} = ft_{1/2} = \ln 2 \times \frac{2\sqrt{2}}{\pi} \frac{f}{f_{FWHM}} \approx 0.62 \frac{f}{f_{FWHM}}$$

$$\frac{z_{1/2}}{\lambda} \approx 0.62 \frac{f}{\delta f_{FWHM}} = 0.62 \frac{\lambda}{\lambda_{FWHM}},$$
 (Gaussian Spectrum)  
$$\frac{z_{1/2}}{\delta f_{FWHM}} = 0.62 \frac{\lambda}{\lambda_{FWHM}},$$

 $\frac{z_{1/2}}{\lambda} \approx 0.62 \times \frac{000}{20} \approx 20$  Wavelengths  $z_{1/2} \approx 12.5 \mu \text{m}$ 

#### Laser vs. LED

• LED (Previous Page: Linewidth = 20nm)

$$rac{z_{1/2}}{\lambda}pprox$$
 0.62  $imes rac{635}{20}pprox$  20Wavelengths  $z_{1/2}pprox$  12.5 $\mu$ m

• HeNe Laser (30cm Coherence Length)

$$\frac{30 \times 10^{-2} \mathrm{m}}{633 \times 10^{-9} \mathrm{m}} \approx 4.7 \times 10^5 \approx 0.62 \frac{f}{f_{FWHM}} = 0.62 \frac{\lambda}{\lambda_{FWHM}}$$

$$f_{FWHM} = \frac{0.62f}{4.7 \times 10^5} = \frac{0.62c}{47 \times 10^6 \lambda} = 620 \times 10^6 \text{Hz}$$

$$\lambda_{FWHM} = \frac{0.62\lambda}{47 \times 10^5} = 830 \times 10^{-15} \text{m}$$

- Linewidth 620MHz or 830fm.

#### Coherence in Beamsplitters

• Beamsplitter 5mm Thick (n = 1.5: OPD 75mm One Way)

$$\frac{z_{1/2}}{\lambda} = \frac{15 \times 10^{-3} \text{m}}{633 \times 10^{-9} \text{m}} = 23700 \approx 0.62 \frac{f}{f_{FWHM}} = 0.62 \frac{\lambda}{\lambda_{FWHM}}$$
$$\lambda_{FWHM} = \frac{0.62 \times 633 \text{nm}}{23700} = 16.6 \text{pm}$$
$$f_{FWHM} = \frac{0.62c}{23700\lambda} = 12.4 \text{GHz}$$

- Laser: Coherent Effects Produce Fringes
- LED:  $z_{1/2}$  Much Less than Thickness: No Fringes

## **Optical Coherence Tomography**

- Short Coherence Source
  - Super–Luminescent
     Diode
  - Ti:Sap Laser
  - Other
- M1 is Reference
- Moving Reference Mirror
- M2 is Target
- Interference? Compare...
  - Path Difference
  - Coherence Length

![](_page_18_Figure_11.jpeg)

## **OCT** Signals

- Examples with Partial Reflectors
- Air-Glass Interfaces (Simulated Signals)
- Idea Extends to Thick "Distributed" Targets

![](_page_19_Figure_4.jpeg)

## **Spatial Coherence**

- Double Slit with Source at A Produces Diffraction Pattern
- Source at B Produces a Different Diffraction Pattern
- What about Both Together?

![](_page_20_Figure_4.jpeg)

## Double Slit with Different Sources

• Sources Mutually Coherent

$$\langle |E_1^* E_2| \rangle = \sqrt{|E_1|^2 |E_2|^2}$$

- Some Well-Defined Pattern with High Visiblity

- Incoherent Sources
  - Each Produces its Own Diffraction Pattern
  - Peaks and Nulls in Different Places
  - Ultimately Incoherent Addition

### Model for Spatial Coherence

- Multiple Point Sources
- Each Independent of the Others
- Compute Coherence Between  $(x_1, y_1, z_1)$ , and  $(x'_1, y'_1, z_1)$

![](_page_22_Figure_4.jpeg)

# Van Cittert–Zernicke Theorem (1)

• Correlation

$$\Gamma(x_1, y_1, x'_1, y'_1) = \left\langle E^*\left(x'_1, y'_1, z_1\right) E\left(x_1, y_1, z_1\right) \right\rangle$$

• Fraunhofer Diffraction

$$E(x_1, y_1, z_1) = \frac{jke^{jkz_1}}{2\pi z_1} e^{jk\frac{x_1^2 + y_1^2}{2z_1}} \int \int E(x, y, 0) e^{jk\frac{x_1^2 + y_1^2}{2z_1}} e^{-jk\frac{xx_1 + yy_1}{z_1}} dxdy$$

$$E^*\left(x_1', y_1', z_1\right) = \frac{-jke^{-jkz_1}}{2\pi z_1}e^{-jk\frac{\left(x_1'\right)^2 + \left(y_1'\right)^2}{2z_1}} \times$$

$$\int \int E^* \left( x', y', 0 \right) e^{-jk \frac{\left( x' \right)^2 + \left( y' \right)^2}{2z_1}} e^{+jk \frac{x' x_1' + y' y_1'}{z_1}} dx' dy'$$

# Van Cittert–Zernicke Theorem (2)

• Combine Previous Equations

$$\Gamma(x_1, y_1, x_1', y_1') = \frac{-jke^{-jkz_1}}{2\pi z_1} e^{-jk\frac{(x_1')^2 + (y_1')^2}{2z_1}} \frac{jke^{jkz_1}}{2\pi z_1} e^{jk\frac{x_1^2 + y_1^2}{2z_1}} \times \\ \int \int \left[ \int \int \left\langle E^*\left(x', y', 0\right) E\left(x, y, 0\right) \right\rangle \times \\ e^{-jk\frac{(x')^2 + (y')^2}{2z_1}} e^{+jk\frac{x^2 + y^2}{2z_1}} e^{+jk\frac{x'x_1' + y'y_1'}{z_1}} e^{-jk\frac{xx_1 + yy_1}{z_1}} dx'dy' \right] dxdy$$

• Simplification:  $x^2 - (x')^2 = (x + x')(x - x')$  $\Gamma(x_1, y_1, x'_1, y'_1) = \frac{-jke^{-jkz_1}}{2\pi z_1} e^{jk \frac{(x'_1)^2 + (y'_1)^2}{2z_1}} \frac{jke^{jkz_1}}{2\pi z_1} e^{-jk \frac{x_1^2 + y_1^2}{2z_1}} \int \int \times \left[ \int \int \langle E^*(x', y', 0) E(x, y, 0) \rangle e^{jk \frac{(x'+x)(x'-x) + (y'+y)(y'-y)}{2z_1}} e^{jk \frac{(xx_1 - x'x'_1) + (yy_1 - y'y'_1)}{z_1}} dx' dy' \right] dxdy$ 

## Van Cittert–Zernicke Theorem (3)

• Random Phases:  $\langle E \rangle = 0$ 

$$\left\langle E^*\left(x',y',0\right)E\left(x,y,0\right)\right\rangle = 0$$

• Analogy with Wiener–Khintchine Theorem Earlier  $\left\langle E^*\left(x',y',0\right)E\left(x,y,0\right)\right\rangle dx'dy' = I\left(x,y,0\right)\delta\left(x-x'\right)\delta\left(y-y'\right)dx'dy'$ 

$$\Gamma(x_1, y_1, x_1', y_1') = \frac{-jke^{-jkz_1}}{2\pi z_1} e^{jk \frac{(x_1')^2 + (y_1')^2}{2z_1}} \frac{jke^{jkz_1}}{2\pi z_1} e^{-jk \frac{x_1^2 + y_1^2}{2z_1}} \int \int \times \left[ \int \int I(x, y, 0) \,\delta\left(x - x'\right) \delta\left(y - y'\right) e^{jk \frac{(x_1 - x'x_1') + (yy_1 - y'y_1')}{z_1}} dx' dy' \right] dx dy$$

# Van Cittert–Zernicke Theorem (4)

• Outer Integral

$$\Gamma(x_1, y_1, x_1', y_1') = \frac{-jke^{-jkz_1}}{2\pi z_1} e^{jk\frac{(x_1')^2 + (y_1')^2}{2z_1}} \frac{jke^{jkz_1}}{2\pi z_1} e^{-jk\frac{x_1^2 + y_1^2}{2z_1}} \times \int \int I(x, y, 0) e^{jk\frac{x(x_1 - x_1') + y(y_1 - y_1')}{z_1}} dxdy$$

• Simplify

$$\Gamma(x_1, y_1, x'_1, y'_1) = \left(\frac{k}{2\pi z_1}\right)^2 e^{jk \frac{(x'_1)^2 + (y'_1)^2}{2z_1}} e^{-jk \frac{x_1^2 + y_1^2}{2z_1}} \times \int \int I(x, y, 0) e^{jk \frac{x(x_1 - x'_1) + y(y_1 - y'_1)}{z_1}} dxdy$$

# Van Cittert–Zernicke Theorem (5)

• Define Coordinate Differences

$$\xi = x_1' - x_1$$
  $\eta = y_1' - y_1$ 

• Van Cittert-Zernicke Theorem

$$\Gamma(\xi,\eta) = \left(\frac{k}{2\pi z_1}\right)^2 \int \int I(x,y,0) e^{jk\frac{x\xi+y\eta}{z_1}} dxdy$$

• Equation Similar to Fraunhofer Diffraction...

$$E(x_1, y_1, z_1) = \frac{jke^{jkz_1}}{2\pi z_1} \int \int E(x, y, 0) e^{-jk\frac{(xx_1 + yy_1)}{z_1}} dxdy$$

• ... but New Meaning (Coherence rather than Field)

## Spatially Coherent and Incoherent Illumination

$$\begin{array}{rccc} E\left(x,y,0\right) & \to & \mathsf{Fraunhofer} & \to & E\left(x_1,y_1,z_1\right) \\ I\left(x,y,0\right) & \to & \mathsf{vC-Z} & \to & \mathsf{\Gamma}\left(x_1-x_1',y_1-y_1'\right) \ \mathsf{at} \ z_1 \end{array}$$

![](_page_28_Picture_2.jpeg)

A. Spatially Coherent (Laser Source Direct)

![](_page_28_Picture_4.jpeg)

B. Spatially Incoherent (Same Size Laser Source on Ground Glass)

## Coherent and Incoherent Source: Calculation

- Coherent Source (1mm HeNe Viewed at 5m)
  - Beam Size d at Receiver
- Spatially Incoherent Source (Same Laser on Ground Glass)
  - Whole Field Illuminated
  - Bright and Dark Regions
  - Diameter of  $\Gamma$  is d
  - "Typical Spot Size:" d

• Spatially and Temporally Incoherent

Source: Filtered Tungsten 400 to 800nm

- Whole Field Illuminated
- Bright and Dark Regions
   Only for Short Times

$$\frac{\delta\lambda}{\lambda} = \frac{\delta f}{f} \to 1$$

 Uniform when Measured for Longer than a Cycle (femtoseconds)

$$d = \frac{4}{\pi} \frac{\lambda}{D} z = \frac{4}{\pi} \frac{633 \times 10^{-9} \text{m}}{10^{-3} \text{m}} \text{ 5m} = 4 \times 10^{-3} \text{m}$$

### Speckle in Scattered Light

- Coherent Source
  - Focused or Diverging
  - Spatially Coherent
  - Temporally Coherent
- Scattering Particles
  - Secondary Sources
  - Randomly Distributed
  - Illuminated by Source

![](_page_30_Picture_9.jpeg)

- Large Illumination Pattern: Small Speckles
- Small Illumination Pattern: Large Speckles
- Infer Size of Illumination by Size of Speckles?
  - See Next Page

#### Example: Finding the Focus

• Example (NA = 0.25): Waist

![](_page_31_Figure_2.jpeg)

$$d_0 = \frac{2}{\pi} \frac{532 \times 10^{-9} m}{0.25} = 1.4 \times 10^{-6} m$$

![](_page_31_Figure_4.jpeg)

• At Glass

![](_page_31_Figure_6.jpeg)

• Speckle Size

$$d_{speckle} = \frac{4}{\pi} \frac{\lambda}{d_{glass}} z$$

- Same as Field of Coherent Source of Size  $d_{glass}$ 

![](_page_31_Figure_10.jpeg)

## Example: Speckle in Laser Radar

- Carbon Dioxide Laser
  - $-\lambda = 10.59 \mu m$
  - $d_0 = 30 \text{cm}$
  - Rayleigh Range

$$b = \pi d_0^2 / (4\lambda) = 6.7$$
 km

• Diameter at Target

$$d_{target} = d_0 \sqrt{1 + \frac{z_t^2}{b^2}}$$

- Far Field  $z_t = 20$ km

 $d_{target} \approx 90 \mathrm{cm}$ 

• Speckle Size at Receiver

$$rac{4}{\pi}rac{\lambda}{d_{target}}z_r=15 imes10^{-6}$$
rad $imes z_r$ 

![](_page_32_Figure_12.jpeg)

• Coaxial System  $z_r = z_t$  and Target in Far–Field

![](_page_32_Figure_14.jpeg)

#### Speckle Decorrelation

- Transverse Motion
  - New Scatterers
  - New Speckle Pattern
- Time on Target (ToT)

$$t_{target} = \frac{d_{target}}{v_{\perp}}$$

$$B_{tot} = \frac{v_{\perp}}{d_{target}}$$

• Recall Doppler Shift

$$f_{DR} = \frac{2v_{\parallel}}{\lambda}$$

![](_page_33_Picture_9.jpeg)

## Speckle Reduction

- Average *n* Measurements
  - Improvement with  $1/\sqrt{n}$
  - Stationary Process
- Temporal Averaging

$$\sigma = \sqrt{\frac{t_{target}}{T_{integration}}}$$

- Long Time
  - Reduced Fluctuation
     but . . .
  - . . . More TargetVariation
- "Matched Filtering"

- Area Averaging
  - Incoherent Averaging
  - Space-Invariant Process
- Aperture Averaging

$$\sigma = \sqrt{\frac{A_{tx}}{A_{rx}}} = \frac{D_{tx}}{D_{rx}}$$

• Coherent Averaging Not Effective

$$\langle E \rangle = 0$$

- Other Diversity
  - Wavelength
  - Polarization

## Making Light Coherent

![](_page_35_Picture_1.jpeg)

![](_page_35_Picture_2.jpeg)

- Interference Filter
  - Dielectric Stack
  - Narrow Bandpass
  - Temporal Coherence

$$\frac{\lambda}{\delta\lambda} = \frac{\delta z}{\lambda}$$

- Pinhole
  - Unresolved
  - Spatial Filter
  - Narrow Spatial
     Bandpass
  - Spatial Coherence

## Making Light Incoherent

- Ground Glass Reduces Spatial Coherence
- Moving it Reduces Temporal Coherence

![](_page_36_Figure_3.jpeg)